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**A PRODUCT OPERATOR THEORY
OF ^{13}C SPIN-ECHO J-MODULATION
NMR SPECTROSCOPY FOR
 CD_n ($n = 1,2,3$) GROUPS**

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ABSTRACT

Product operator formalism is widely used for analytical description of multiple-pulse NMR experiments for a weakly coupled spin systems. ^{13}C spin-echo J-modulation NMR spectroscopy for CH_n ve CD_n groups is used for identification of different carbon groups. In this study, by using the product operator technique, the analytical description of ^{13}C spin-echo J-modulation NMR spectroscopy for CD_n ($n = 1,2,3$) groups is presented and the experimental identifications of ^{13}C NMR signals of CD_3 , CD_2 and CD groups and also quaternary carbons are discussed.

INTRODUCTION

In order to analyze multiple-pulse NMR experiments applied to large spin systems, a quantum mechanical approach has to be used. In analytical description of multiple-pulse 1D and 2D NMR experiments, the product operator formalism is widely used for weakly coupled spin systems¹⁻⁷. Spin-echo J-modulation is a well known and widely applied technique, *e.g.* for signal assignments in ¹³C NMR spectroscopy and for 2D ¹³C NMR spectroscopy^{8,9}. So, by using spin-echo J-modulation, ¹³C NMR signals coming from CH_n (n=0,1,2,3) groups can be assigned¹⁰. It also can be used for the elimination of ¹³C NMR signals from deuterated solvents and for the ¹³C NMR analysis of partially deuterated compounds¹¹. The product operator description of heteronuclear 2D J-resolved NMR and 2D DEPT J-resolved NMR spectroscopies for the weakly coupled IS_n spin systems (I=1/2, S=1) has been reported elsewhere^{12,13}. In this study, the product operator technique is used for analytical description of ¹³C spin-echo J-modulation NMR spectroscopy for CD_n (n=1,2,3) groups by representing CD_n as IS_n (n=1,2,3, I=1/2, S=1). Then the experimental identifications of ¹³C NMR signals coming from CD₃, CD₂, CD and quaternary carbons are discussed. To the best of our knowledge, this will be the first application of product operator theory to ¹³C spin-echo J-modulation NMR spectroscopy for CD_n (n=1,2,3) groups.

THEORY

Time dependency of the density matrix is given by⁷

$$\sigma(t) = \exp(-iHt)\sigma(0)\exp(iHt) \quad (1)$$

Where *H* is the total Hamiltonian which consists of radio frequency (r.f.) pulse, chemical shift and spin-spin coupling Hamiltonians and $\sigma(0)$ is the density matrix at *t*=0. After employing the Hausdorff formula⁷

$$\begin{aligned} \exp(-iHt)A\exp(iHt) &= A - (it)[H, A] + \frac{(it)^2}{2!}[H, [H, A]] \\ &\quad - \frac{(it)^3}{3!}[H, [H, [H, A]]] + \dots, \end{aligned} \quad (2)$$

the r.f pulse, chemical shift and spin-spin coupling evolution of product operators can easily be obtained^{1,6,7}. The details on the evolution of product operators under these Hamiltonians can be found elsewhere^{1,6,7,12}. At any time during the experiment, the ensemble averaged expectation value of the spin angular momentum, *e.g.* for I_y , is

$$\langle I_y \rangle = \text{Tr}(I_y \sigma(t)). \quad (3)$$

Where $\sigma(t)$ is the density matrix operator calculated from Eq. (1) at any time. As $\langle I_y \rangle$ is proportional to the magnitude of the y magnetization, it represents the signal detected on y axis. So, in order to estimate the FID signal of a multiple-pulse NMR experiment, density matrix operator should be obtained at the end of the experiment.

^{13}C Spin-Echo J-Modulation NMR For CD_n ($n = 1, 2, 3$)

When the simple pulse sequence of spin-echo J-modulation is used, there exists some sources of errors¹⁴. In order to eliminate these sources of errors, a new pulse sequence called ESCORT editing has been developed¹⁴. In this study, for the analytical description of ^{13}C spin-echo J-modulation NMR spectroscopy of CD_n (IS_n) groups, the pulse sequence of ESCORT editing, illustrated in Figure 1, is used. In the pulse sequence τ_0 takes an average value found from average coupling constant J_0 as $\tau_0 = 1/J_0$ and it is kept constant. As seen in Figure 1, the density matrix operator at each stage of the experiment is labeled with numbers.

For **IS** spin system s_I is the density matrix operator at thermal equilibrium and

$$\sigma_1 = I_z. \quad (4)$$

The pulse sequence obviously leads to the following density matrices for each labeled point:

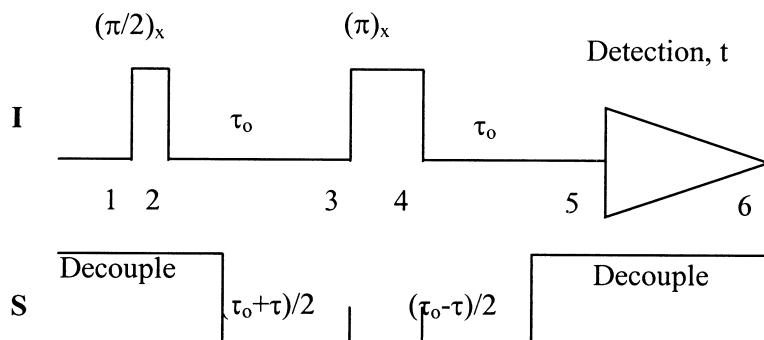


Figure 1. The pulse sequence of ESCORT editing for spin-echo J-modulation NMR spectroscopy¹⁴.

$$\sigma_2 = -I_y \quad (5)$$

$$\sigma_3 = I_x S_z S_{J0} - I_y (1 + S_z^2 (C_{J0} - 1)) \quad (6)$$

$$\sigma_4 = I_x S_z S_{J0} + I_y (1 + S_z^2 (C_{J0} - 1)). \quad (7)$$

In the preceding and following equations $S_{J0} = \sin \pi J(\tau_0 + \tau)$, $C_{J0} = \cos \pi J(\tau_0 + \tau)$, $C_{nJ} = \cos 2\pi n J\tau$ and $C_I = \cos \Omega_I t$. It is assumed that during τ_0 , relaxation and evolution under chemical shift do not exist. The chemical shift evolution take place only during t_1 . Then,

$$\sigma_5 = I_y (1 + S_z^2 (C_J - 1)) \quad (8)$$

and

$$\sigma_6 = I_y (1 + S_z^2 (C_J - 1)) C_I \quad (9)$$

As seen in Figure 1, s_6 is the density matrix operator at the end of the experiment. For **IS** spin system in s_6 only the observable I_y and $I_y S_z^2$ terms are kept. Now, it is necessary to obtain the $Tr(I_y O)$ values of observable product operators indicated by O . For IS_n spin system ($I = 1/2$, $S = 1$, $n = 1, 2, 3$), $Tr(I_y O)$ values for some of the observable product operators were calculated by a computer program and the results are given in Table 1. So using Table 1

$$\langle I_y \rangle = Tr(I_y \sigma_6) = \left(\frac{1}{2} + C_J \right) C_I \quad (10)$$

is obtained.

By using the same pulse sequence for **IS₂ spin system**, we obtain

$$\sigma_6 = I_y (1 + (S_{1z}^2 + S_{2z}^2) (C_J - 1) + S_{1z}^2 S_{2z}^2 (C_J - 1)^2) C_I. \quad (11)$$

Table 1. The Results of the $Tr(I_y O)$ Calculations for Some of the Observable Product Operators in IS_n Spin System ($I = 1/2$, $S = 1$, $n = 1, 2, 3$)

Spin System	Product Operator (O)	$Tr(I_y O)$
<i>IS</i>	I_y	3/2
	$I_y S_z^2$	1
<i>IS₂</i>	I_y	9/2
	$I_y (S_{1z}^2 + S_{2z}^2)$	6
	$I_y S_{1z}^2 S_{2z}^2$	2
<i>IS₃</i>	I_y	27/2
	$I_y (S_{1z}^2 + S_{2z}^2 + S_{3z}^2)$	27
	$I_y (S_{1z}^2 S_{2z}^2 + S_{1z}^2 S_{3z}^2 + S_{2z}^2 S_{3z}^2)$	18
	$I_y S_{1z}^2 S_{2z}^2 S_{3z}^2$	4

Then,

$$Tr(I_y\sigma_6) = \left(\frac{3}{2} + 2C_J + C_{2J}\right)C_I \quad (12)$$

is obtained.

For **IS₃ spin system**, applying the same procedure we obtain

$$\sigma_6 = I_y \left(\begin{array}{l} 1 + (S_{1z}^2 + S_{2z}^2 + S_{3z}^2)(C_J - 1) + (S_{1z}^2 S_{2z}^2 + S_{1z}^2 S_{3z}^2 + S_{2z}^2 S_{3z}^2)(C_J - 1)^2 \\ + S_{1z}^2 S_{2z}^2 S_{3z}^2 (C_J - 1)^3 \end{array} \right) C_I \quad (13)$$

and

$$Tr(I_y\sigma_6) = \left(\frac{7}{2} + 6C_J + 3C_{2J} + C_{3J}\right)C_I \quad (14)$$

DISCUSSION

$Tr(I_y\sigma_6)$ values obtained in the preceding section for IS, IS₂ and IS₃ spin systems represents the FID signals of ¹³C spin-echo J-modulation NMR spectroscopy for CD, CD₂ and CD₃ groups, respectively. $Tr(I_y\sigma_6)$ values can be normalized multiplying by 4/ $Tr(E)$. Here E is the unity product operator for the corresponding spin system. Then the normalized values become as following:

$$Tr(I_y\sigma_6)(IS) = \frac{2}{3} \left(\frac{1}{2} + C_J\right) C_I \quad (15)$$

$$Tr(I_y\sigma_6)(IS_2) = \frac{2}{9} \left(\frac{3}{2} + 2C_J + C_{2J}\right) C_I \quad (16)$$

$$Tr(I_y\sigma_6) = \frac{2}{27} \left(\frac{7}{2} + 6C_J + 3C_{2J} + C_{3J}\right) C_I \quad (17)$$

By adding the $Tr(I_y\sigma_6)$ values of all these three groups, total $Tr(I_y\sigma_6)$ value can be found as:

$$(Tr(I_y\sigma_6))_{\text{tot}} = \left(\begin{array}{l} \frac{2}{3} \left(\frac{1}{2} + C_J\right) + \frac{2}{9} \left(\frac{3}{2} + 2C_J + C_{2J}\right) \\ + \frac{2}{27} \left(\frac{7}{2} + 6C_J + 3C_{2J} + C_{3J}\right) \end{array} \right) C_I. \quad (18)$$

This corresponds to total FID signal for all three groups. This does not include the FID signal of quaternary carbons. By taking $\theta = 2\pi J_0\tau$, FID signals can be estimated for different angles (for different τ values). When the experiment performed at 120° or 240° angles, only the signals coming from quaternary carbons are observed as $(Tr(I_y\sigma_6))_{tot}$ value becomes zero at these angles. For identification of ^{13}C NMR signals of CD, CD_2 and CD_3 groups, the FID signals of ^{13}C spin-echo J-modulation NMR spectroscopy should be obtained at 90° and 180° angles (at corresponding τ values). Then, from their combinations, following results are obtained:

$$I(\text{CD}_2) = 4.5\{(Tr(I_y\sigma_6))_{tot}(90^\circ) + (Tr(I_y\sigma_6))_{tot}(180^\circ)\} \quad (19)$$

$$I(\text{CD}) + \frac{1}{9}I(\text{CD}_3) = 1.5\{(Tr(I_y\sigma_6))_{tot}(90^\circ) - Tr(I_y\sigma_6))_{tot}(180^\circ)\} \quad (20)$$

where $I(\text{CD}_n)$ represents the relative signal intensity of indicated group. The combination in Equation (19) gives the FID signal of CD_2 group only. And the combination in Equation (20) results in the FID signals of CD and CD_3 groups. It is obvious from Equation (20) that CD signals can be easily separated from CD_3 signals as CD_3 signals will have small relative intensities.

In conclusion, product operator theory is applied to ^{13}C spin-echo J-modulation NMR spectroscopy for CD_n ($n=1,2,3$) groups. In order to identify the CD_3 , CD_2 and CD groups and also quaternary carbons, experimental suggestions for ^{13}C spin-echo J-modulation NMR spectroscopy of deuterated molecules were made.

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